

# IMBq 方程的多辛格式及其守恒律

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摘要 考虑非线性 IMBq 方程的多辛 Hamilton 形式, 通过消去中间变量, 得到新的等价于多辛 Preissman 积分的格式. 发现它具有多辛守恒律、局部能量守恒律及局部动量守恒律, 最后以数值例子验证其有效性.

关键词 IMBq 方程, 多辛, 守恒律, 非线性, Hamilton 形式

中图分类号 O 241. 82

文献标识码 A

由冯康先生开创的 Hamilton 系统的辛几何算法<sup>[1]</sup>在理论上已日趋成熟, 且被广泛地应用于大量的数学物理方程. 这一算法因其独特的稳定性与长期跟踪能力已解决了许多实际问题<sup>[2, 3]</sup>, 这是传统非辛算法所不能比拟的. 因此, 对辛算法的研究具有重要的理论与实际意义. 本文的目的在于用辛几何的观点, 解决下列非线性 IMBq 方程的周期初值问题

$$\left. \begin{aligned} u_t - u_{xxt} &= (u^3)_{xx} + u_{xx}, \\ u(x+80, t) &= u(x, t), \\ u(x, 0) &= f(x). \end{aligned} \right\} \quad (1)$$

它用于描述一定限制条件下弹性杆的纵向形变波的传播, 在离子声波方面应用较为广泛. 本文就用多辛算法研究此方程.

## 1 非线性 IMBq 方程的多辛形式

Reich 曾指出大量偏微分方程都具有多辛结构<sup>[4, 5]</sup>, 从而可改写多辛 Hamilton 方程组形式为

$$Mz_t + Kz_x = -zS(z), \quad z \in \mathbf{R}^n, \quad (x, t) \in \mathbf{R}^2. \quad (2)$$

在式(2)中,  $M, K \in \mathbf{R}^{n \times n} (n \geq 3)$  是反对称矩阵,  $S: \mathbf{R}^n \rightarrow \mathbf{R}$  是光滑函数, 称为 Hamilton 函数,  $-zS(z)$  为函数  $S(z)$  的梯度. 多辛 Hamilton 系统(2)满足下列几个守恒律. ( ) 多辛守恒律. 有

$$\frac{\partial}{\partial t} \omega + \frac{\partial}{\partial x} k = 0, \quad (3)$$

其中  $\omega = \frac{1}{2} dz \wedge M dz, k = \frac{1}{2} dz \wedge K dz$ . ( ) 局部能量守恒律. 有

$$\frac{\partial}{\partial t} E(x, t) + \frac{\partial}{\partial x} F(x, t) = 0, \quad (4) \text{ 其中}$$

$E(x, t) = S(z) - \frac{1}{2} z^T K z_x, F(x, t) = \frac{1}{2} z^T K z_t$ . ( ) 局部动量守恒律. 有

$$\frac{\partial}{\partial t} I(x, t) + \frac{\partial}{\partial x} G(x, t) = 0, \quad (5)$$

其中  $I(x, t) = \frac{1}{2} z^T M z_x, G(x, t) = S(z) - \frac{1}{2} z^T M z_t$ . 引入正则动量  $w_x = v, u_t = v_x, u_x = q, v_t = p$ , 则可得方程组(1)的多辛 Hamilton 方程组

收稿日期 2005-01-07

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基金项目 华侨大学科研基金资助项目(04HZR08)

$$\left. \begin{aligned} w_t - p_x &= V'(u) + u, & -q_t - w_x &= -v_{xx} - v, \\ -u_t + v_x &= 0, & u_x &= q, & v_t &= p. \end{aligned} \right\} \tag{6}$$

其中  $V(u) = u^4/4$ . 令状态变量  $z = (u, v, w, p, q)^T$ , 于是 Hamilton 方程组(6)可写成形式(2). 此时

$$S = \frac{1}{2}(u^2 - v^2 + v_x^2) + pq + V(u),$$
$$M = \begin{bmatrix} 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \end{bmatrix}, \quad K = \begin{bmatrix} 0 & 0 & 0 & -1 & 0 \\ 0 & 0 & -1 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{bmatrix}.$$

2 几个局部守恒律

易证方程组(6)满足上述几个局部守恒律, 其具体形式有 3 种(因篇幅关系证明从略). (i) 相应于式(3)的多辛守恒律为

$$\frac{\partial}{\partial t}(du \wedge dw + dq \wedge dv) + \frac{\partial}{\partial x}(dp \wedge du + dw \wedge dv) = 0. \tag{7}$$

(ii) 相应于式(4)的局部能量守恒律为

$$\frac{\partial}{\partial t}[\frac{1}{2}(u^2 + v_x^2 - v^2) + V(u) - wv_x] + \frac{\partial}{\partial x}(2pu_t + wv_t) = 0. \tag{8}$$

(iii) 相应于式(5)的局部动量守恒律为

$$\frac{\partial}{\partial t}(uw_x + qv_x) + \frac{\partial}{\partial x}[\frac{1}{2}(u^2 - v^2 + u_t^2) + V(u) - uv_t] = 0. \tag{9}$$

3 Preissman 多辛积分及其等价格式

引入正则动量  $\frac{1}{\Delta x}(w_{i+1}^{j+1/2} - w_i^{j+1/2}) = v_{i+1}^{j+1/2}, \frac{1}{\Delta t}(u_{i+1}^{j+1/2} - u_i^{j+1/2}) = \frac{1}{\Delta x}(v_{i+1}^{j+1/2} - v_i^{j+1/2}), p_{i+1}^{j+1/2} = \frac{1}{\Delta t} \cdot (v_{i+1}^{j+1/2} - v_i^{j+1/2}), q_{i+1}^{j+1/2} = \frac{1}{\Delta x}(u_{i+1}^{j+1/2} - u_i^{j+1/2})$ , 则得到方程组(6)的中心 Preissman 格式为

$$\left. \begin{aligned} \frac{1}{\Delta t}(w_{i+1}^{j+1/2} - w_i^{j+1/2}) - \frac{1}{\Delta x}(p_{i+1}^{j+1/2} - p_i^{j+1/2}) &= V'(u_{i+1}^{j+1/2}) + u_{i+1}^{j+1/2}, \\ \frac{1}{\Delta t}(q_{i+1}^{j+1/2} - q_i^{j+1/2}) + \frac{1}{\Delta x}(w_{i+1}^{j+1/2} - w_i^{j+1/2}) &= (v_{i+1}^{j+1/2})_{xx} + v_{i+1}^{j+1/2}, \\ \frac{1}{\Delta t}(u_{i+1}^{j+1/2} - u_i^{j+1/2}) &= \frac{1}{\Delta x}(v_{i+1}^{j+1/2} - v_i^{j+1/2}), \\ \frac{1}{\Delta t}(v_{i+1}^{j+1/2} - v_i^{j+1/2}) &= p_{i+1}^{j+1/2}, \quad \frac{1}{\Delta x}(u_{i+1}^{j+1/2} - u_i^{j+1/2}) = q_{i+1}^{j+1/2}. \end{aligned} \right\} \tag{10}$$

消去中间变量  $v, w, p, q$ , 就得到一个与 Preissman 多辛积分等价的新格式为

$$\frac{\delta_x^2(u_{i+3}^{j+1} + u_{i+3}^j) + 4\delta_x^2(u_{i+2}^{j+1} + u_{i+2}^j) + 6\delta_x^2(u_{i+1}^{j+1} + u_{i+1}^j) + 4\delta_x^2(u_i^{j+1} + u_i^j) + \delta_x^2(u_{i-1}^{j+1} + u_{i-1}^j)}{32\Delta t^2} -$$
$$\frac{\delta_x^2\delta_x^2(u_{i+3}^{j+1} + u_{i+3}^j) - 4\delta_x^2\delta_x^2(u_{i+2}^{j+1} + u_{i+2}^j) - 10\delta_x^2\delta_x^2(u_{i+1}^{j+1} + u_{i+1}^j) - 4\delta_x^2\delta_x^2(u_i^{j+1} + u_i^j) + \delta_x^2\delta_x^2(u_{i-1}^{j+1} + u_{i-1}^j)}{32\Delta t^2\Delta x^2} =$$
$$\frac{1}{8}[V'(u_{i+1}^{j+1}) + V'(u_i^{j+1}) + 2(V'(u_{i+1}^j) + V'(u_i^j)) + V'(u_{i-1}^{j-1}) + V'(u_{i-1}^j)]_{xx} +$$
$$\frac{1}{8}[u_{i+1}^{j+1} + u_i^{j+1}) + 2(u_{i+1}^j + u_i^j) + u_{i+1}^{j-1} + u_i^{j-1}]_{xx}. \tag{11}$$

初始条件的离散化处理从略. 其中  $u_i^j \approx u(i\Delta x, j\Delta t)$ .  $\Delta x$  和  $\Delta t$  分别是空间步长和时间步长, 且  $u_{i+1/2}^j =$

$$\frac{1}{2} \cdot (u_i^j + u_{i+1}^j), \quad u_i^j = u_{i+1/2}^{j+1/2} = \frac{1}{2}(u_{i+1/2}^{j+1} + u_{i+1/2}^{j-1}) = \frac{1}{4}(u_i^j + u_{i+1}^{j+1} + u_{i+1}^{j-1} + u_i^{j-1}), \quad u_i^{j+1/2} = \frac{1}{2}(u_i^{j+1} + u_i^j),$$

$(u_i^j)_{xx} = \frac{\mathfrak{L} u_i^j}{\Delta x^2}, \quad \mathfrak{L} u_i^j = u_{i+1}^j - 2u_i^j + u_{i-1}^j, \quad \mathfrak{L} u_i^j = u_{i+1}^{j+1} - 2u_i^j + u_{i-1}^{j-1} \text{ 等等.}$

注 格式(10)具离散多辛守恒律<sup>[6]</sup>. 有

$$\frac{(du_{i+1/2}^{j+1/2} \wedge dw_{i+1/2}^{j+1/2} - du_{i+1/2}^j \wedge dw_{i+1/2}^j) + (dq_{i+1/2}^{j+1/2} \wedge dv_{i+1/2}^{j+1/2} - dq_{i+1/2}^j \wedge dv_{i+1/2}^j)}{\Delta t} + \frac{(dp_{i+1}^{j+1/2} \wedge du_{i+1}^{j+1/2} - dp_{i+1}^j \wedge du_{i+1}^j) + (dw_{i+1}^{j+1/2} \wedge dv_{i+1}^{j+1/2} - dw_{i+1}^j \wedge dv_{i+1}^j)}{\Delta x} = 0. \quad (12)$$

## 4 数值例子

在非线性 IMBq 方程的周期初值问题(1)中取  $f(x) = A \operatorname{sech} kx$ , 设  $u(-40, t) = u(40, t)$ , 其精确解为  $u(x, t) = A \operatorname{sech}(kx + ct)$ , 其中  $A, k, c$  满足关系式  $A^2 = 2c^2, k^2 = \frac{c^2}{c^2 + 1}$ . 取  $\Delta x = 0.2, \Delta t = 0.001, A = \sqrt{2}, k = 1/\sqrt{2}$ , 用格式(11)进行数值模拟. 为方便起见, 除了  $t = 0$  时的初始层采用精确值外, 第 2, 3 层也用精确值进行计算. 图 1 画出了计算 10 000 层后格式(11)在区间  $[-40, 40]$  上的数值解与原方程的精确解的结果比较情况. 结果表明多辛格式(11)是有效的, 具有良好的长时间数值行为, 数值结果与理论分析相符.

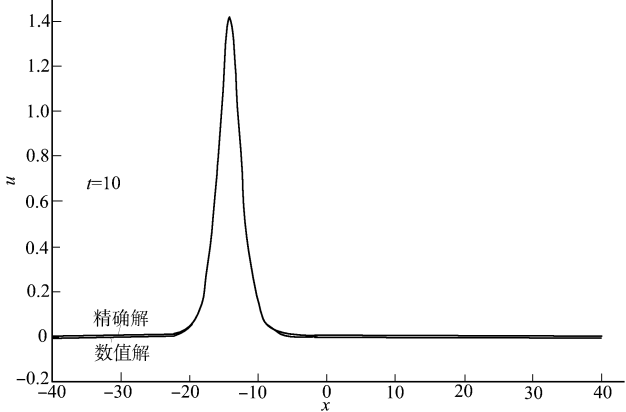


图 1 数值结果比较图

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# Multi-Symplectic Scheme for IMBq Equation and the Conservation Laws to Which It Subjects

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**Abstract** In consideration of multi symplectic Hamiltonian formulation for the nonlinear IMBq equation, by eliminating intermediate variables, the author gets a new scheme which is equivalent to the multi symplectic Preissman integrator. Scheme subjects to conservation laws of multic symplecticity, local energy and local momentum. Finally, its effectiveness is verified by numerical example.

**Keywords** IMBq equation, multi symplecticity, conservation law, nonlinear, Hamiltonian formulation